

Planetary Population Synthesis

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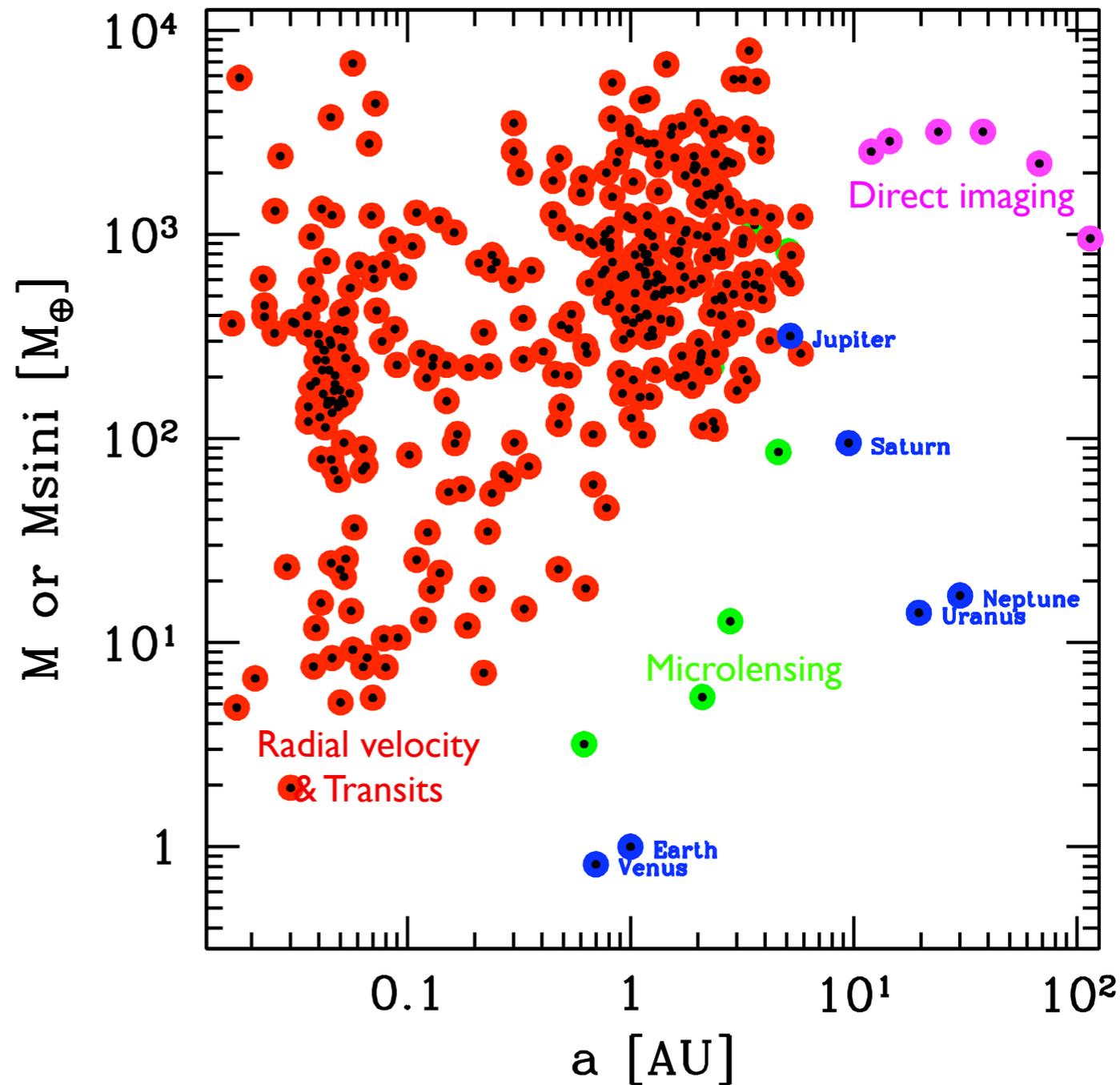
Jackson Hole, 16.9.2011

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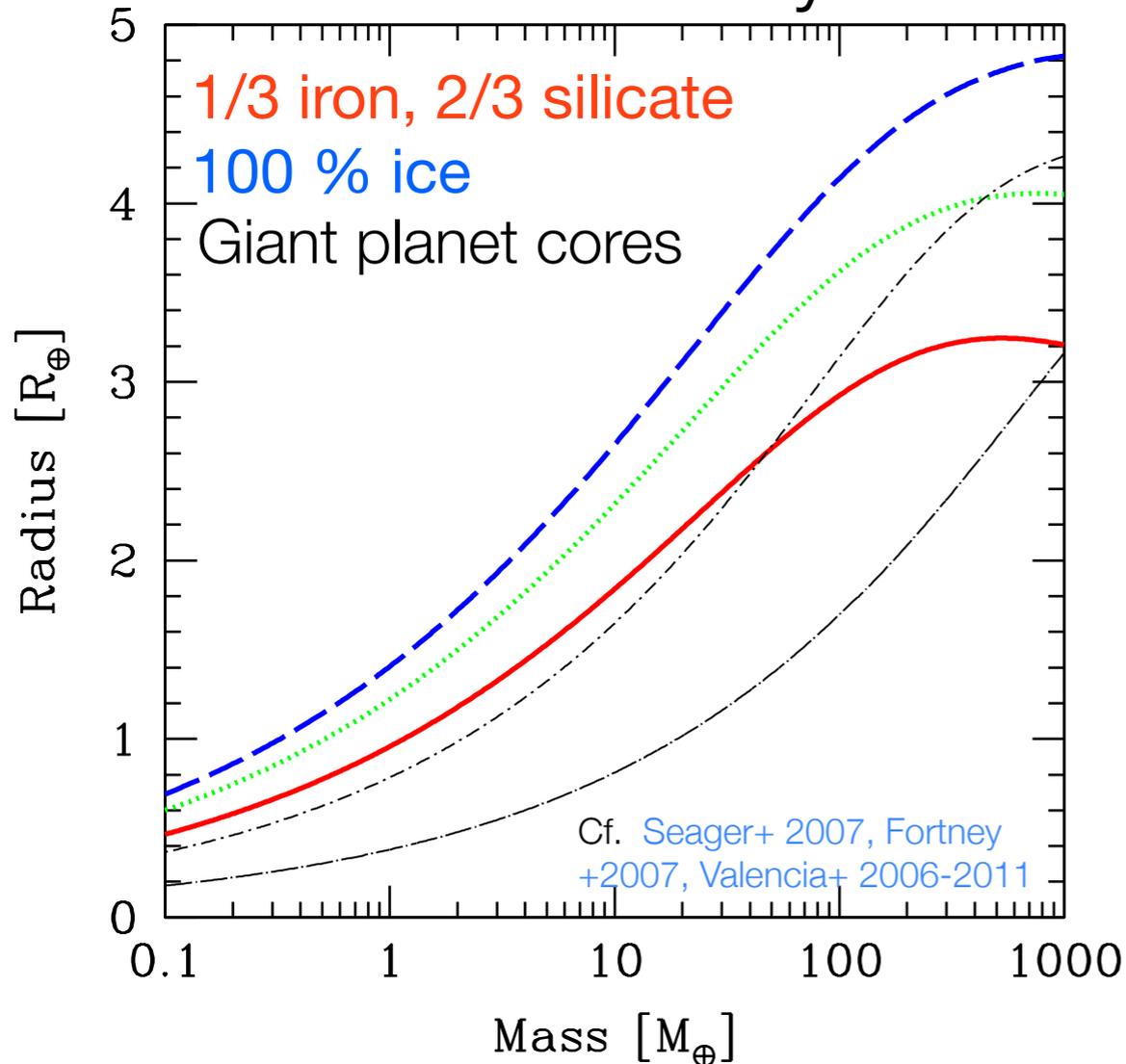
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I Model

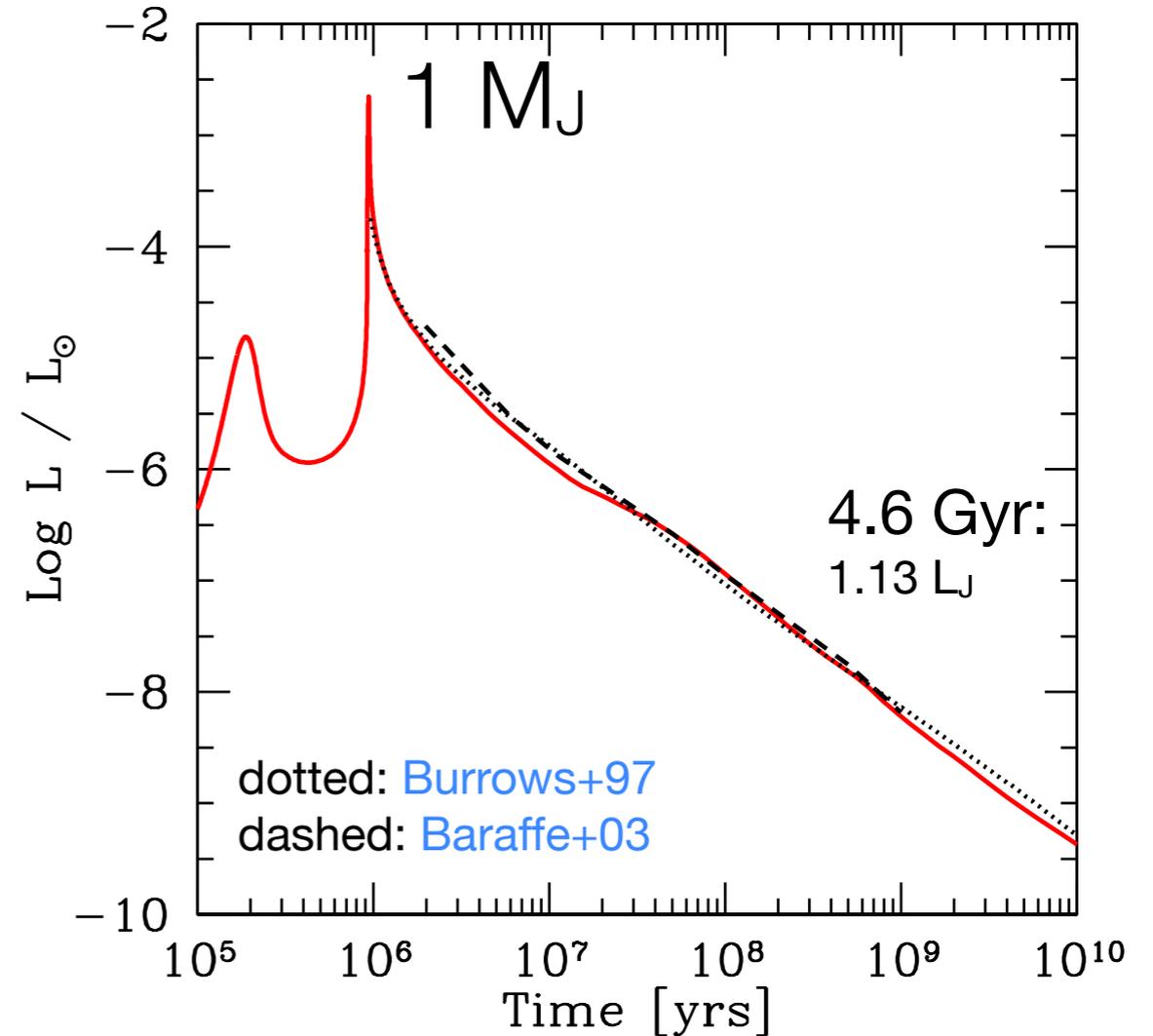
Core accretion + disk evolution + migration

Upgraded planet model

Variable core density



Coupled long term evolution



Simple three layer **differentiated** planet model:

- iron/nickel, silicates, and if accretion at $a > a_{\text{snow}}$ ices. EOS from Seager+ 2007
- Includes effect of **external pressure** (cores).
- Core luminosity from **radioactive** decay, assuming chondritic composition.

-Simple, **grey** atmosphere

$$P = \frac{2g}{3\kappa} \quad T_{\text{int}}^4 = \frac{L_{\text{int}}}{4\pi\sigma R^2}$$

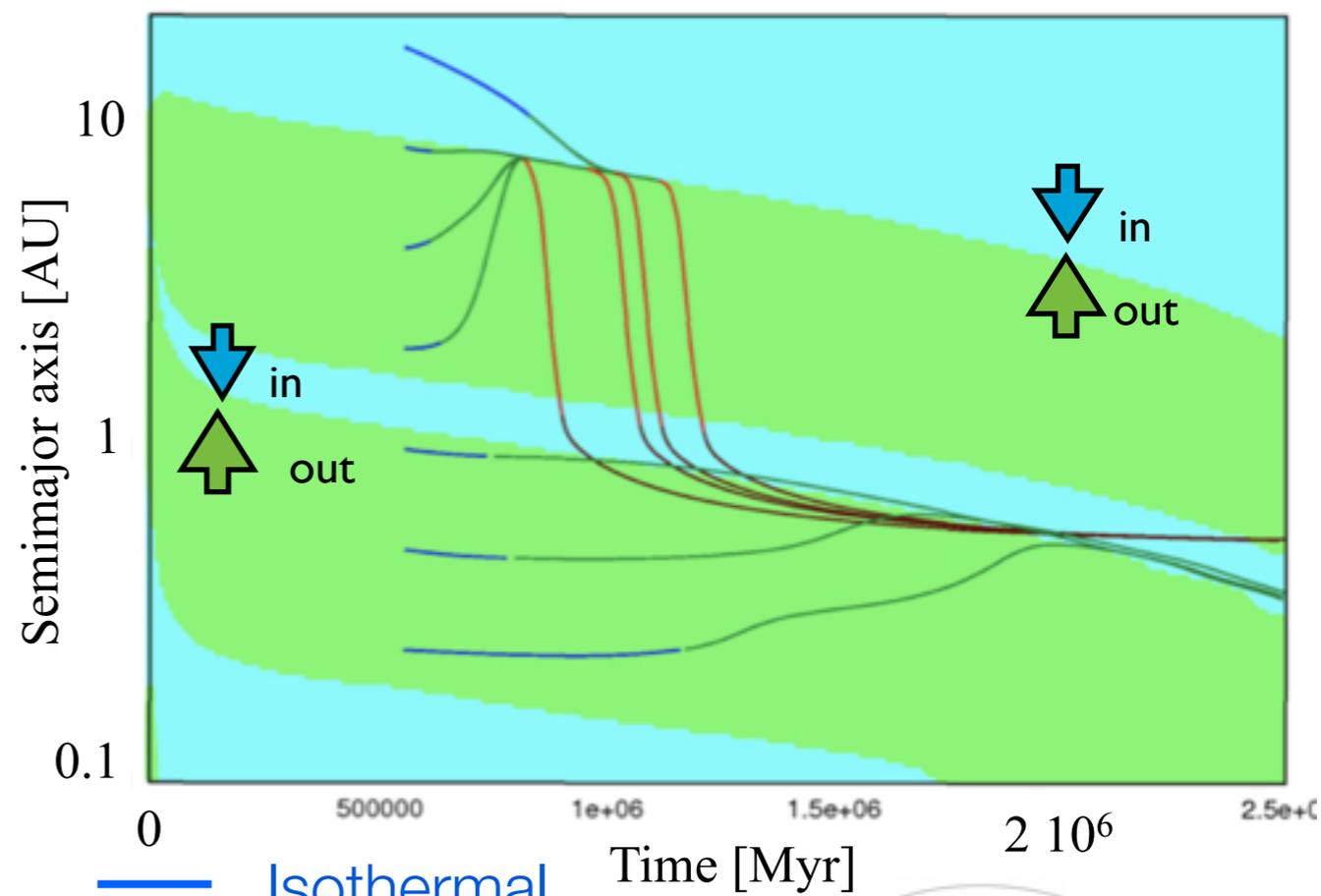
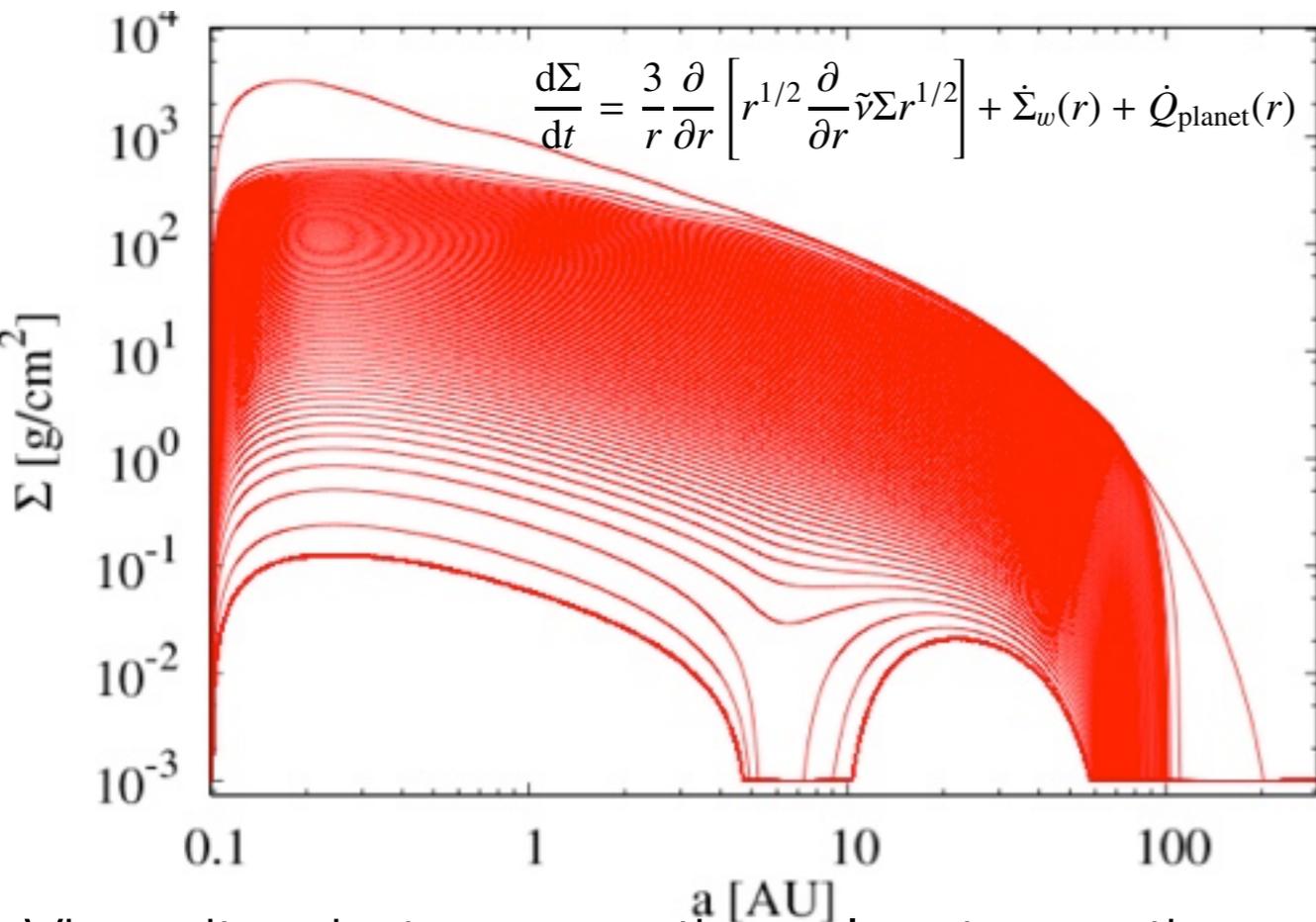
$$T_{\text{equi}} = 280 \text{ K} \left(\frac{a}{\text{IAU}}\right)^{-\frac{1}{2}} \left(\frac{M_*}{M_{\odot}}\right) \quad T^4 = (1 - A)T_{\text{equi}}^4 + T_{\text{int}}^4$$

- Gas **opacities** from Freedman+ 2008

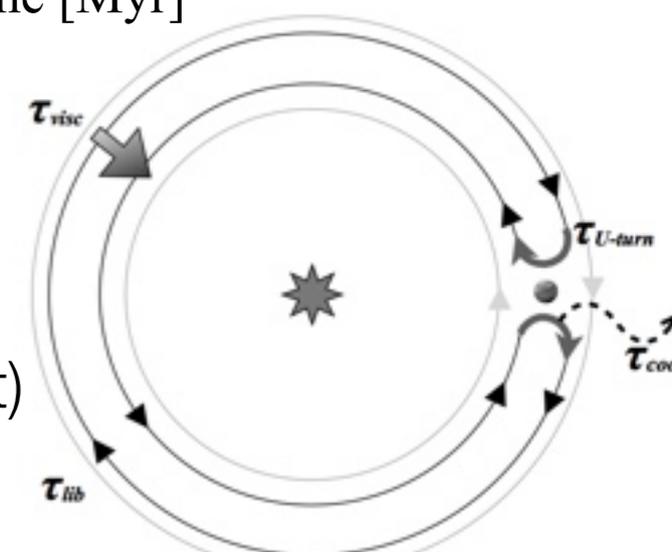
-Gives radius and luminosity at **any** moment, inkl. collapse (Lissauer+ 2009)

Upgraded disk & migration model

Disk gas surface density (1+1D α model) Non-isothermal type I migration



- Isothermal
- Adiabatic
- Ad. saturated
- Type II



Different regimes (in/out) found w. timescales

(Kley & Crida 2008, Kley et al. 2009, Paardekooper et al. 2010, Baruteau & Lin 2010, Bitsch et al. 2011...)

Viscosity, photoevaporation, planet accretion.

New:

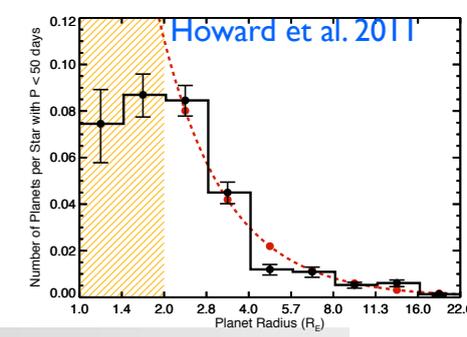
- stellar irradiation
- external photoevaporation [Matsuyama+ 2003](#)
- internal photoevaporation [Clarke+ 2001](#)
- updated initial profile [Andrews+ 2009](#)

$$\Sigma(r, t = 0) = \Sigma_0 \left(\frac{r}{R_0} \right)^{-\gamma} \exp \left[- \left(\frac{r}{R_c} \right)^{2-\gamma} \right]$$

II Population synthesis

Vary initial conditions of formation model according to observed distributions of metallicity, disk mass and disk lifetime.

Formation tracks - Kepler comparison



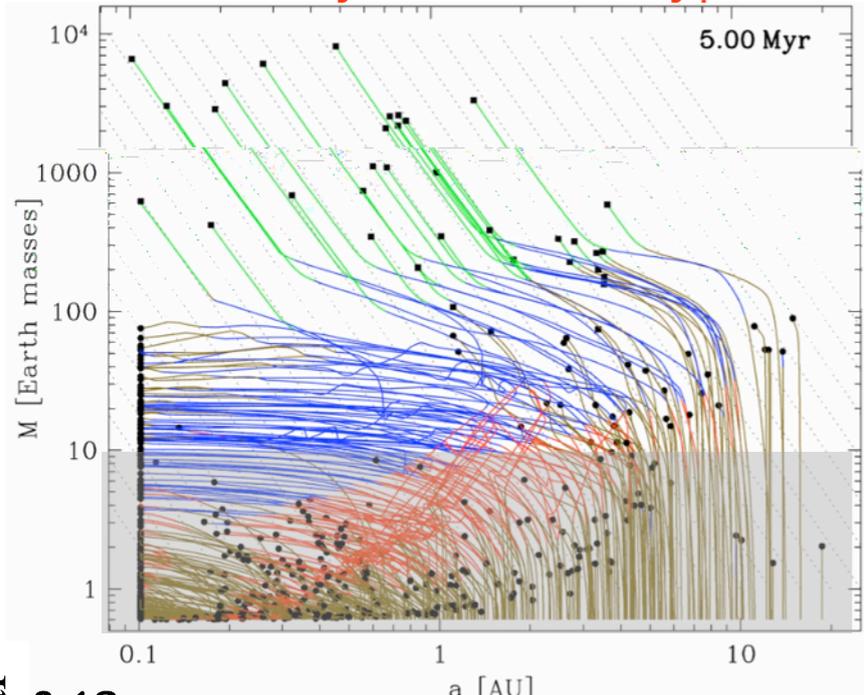
$M_{\text{star}} = 1 M_{\text{sun}}$, $M_{\text{emb},0} = 0.6 M_{\text{earth}}$, irradiated disk, viscosity $\alpha = 7 \times 10^{-3}$

Nominal Model.
Updated Type I.

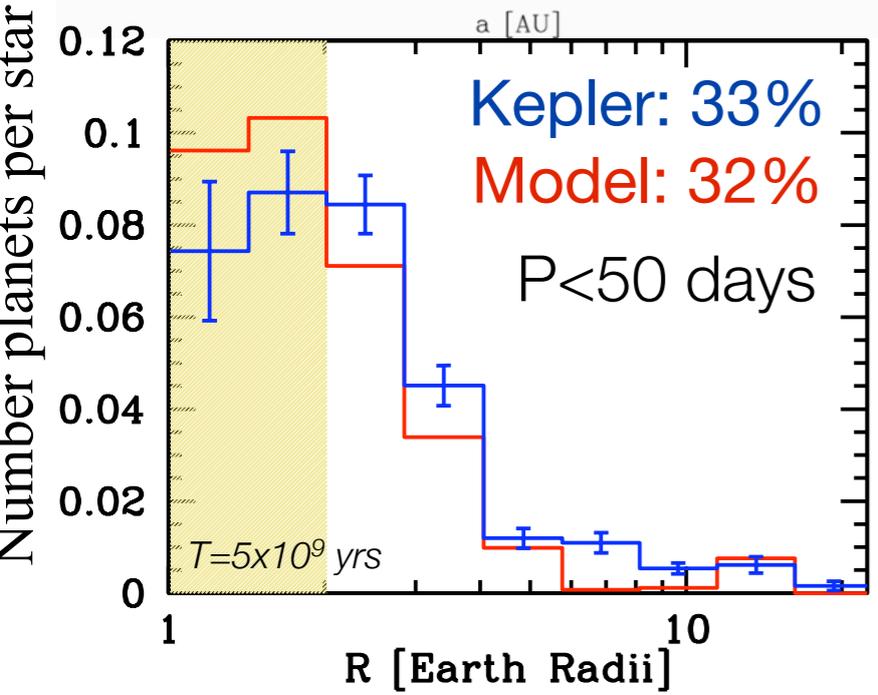
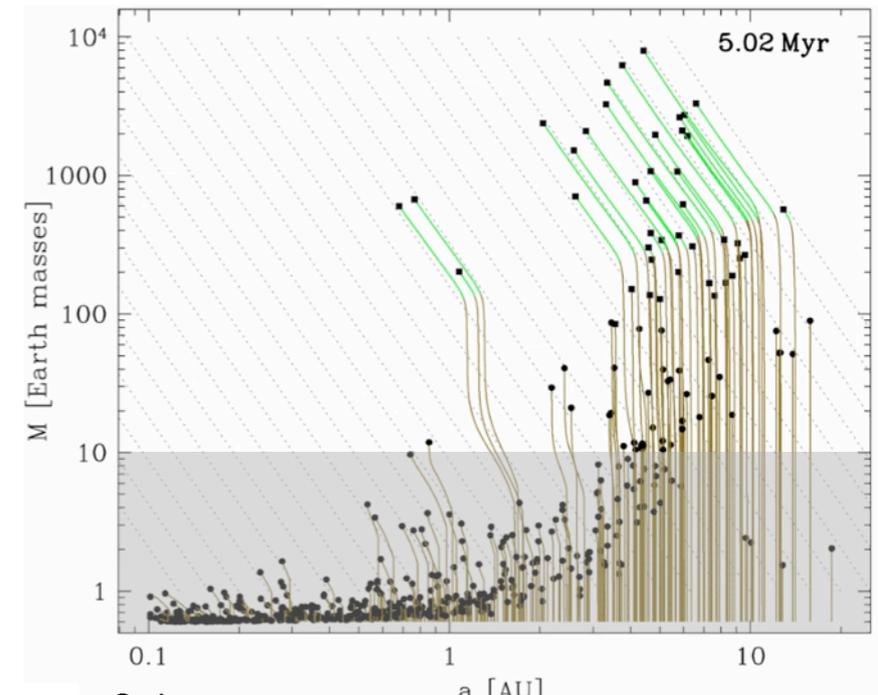
Isothermal Type I.
Tanaka $f_1 = 0.01$

Almost no type I migration.

No efficiency factor for type I!

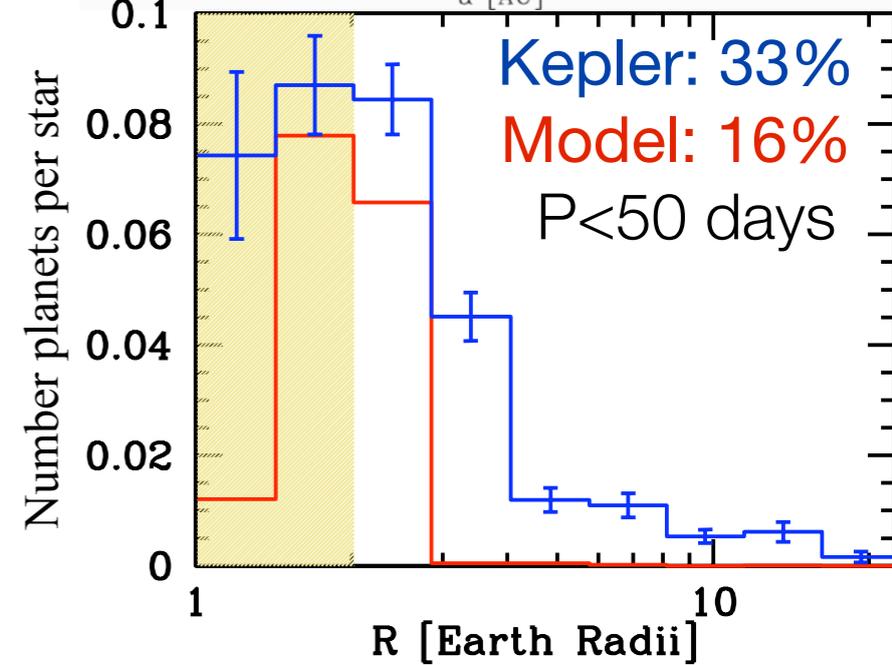


- Diversity!
- Can form both giant planets and close-in low mass planets.
- Migration still too rapid (giants).
- Imprint of convergence zones.



- Generally good agreement:
 - Increase towards small radii.
 - Many low radii.
 - Hot Jupiter 0.5-1%
 - Even in absolute fraction.

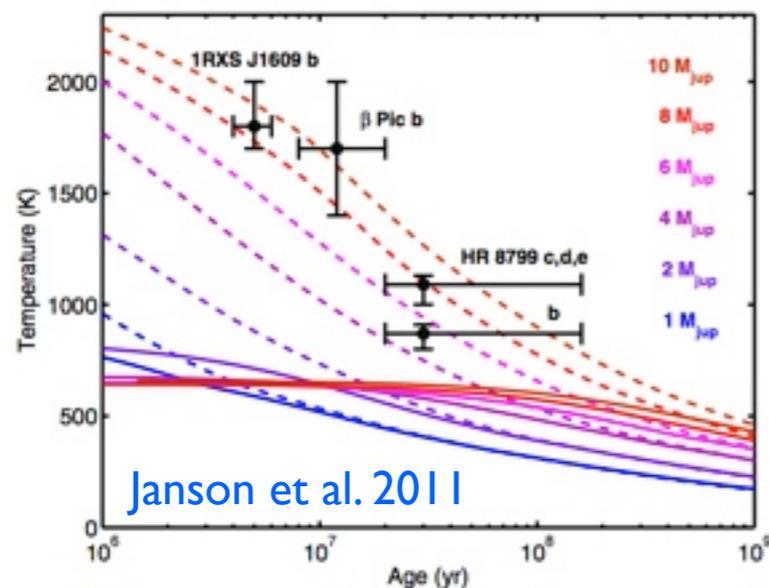
But no bimodality...!



Luminosity of young Jupiters *revisited*

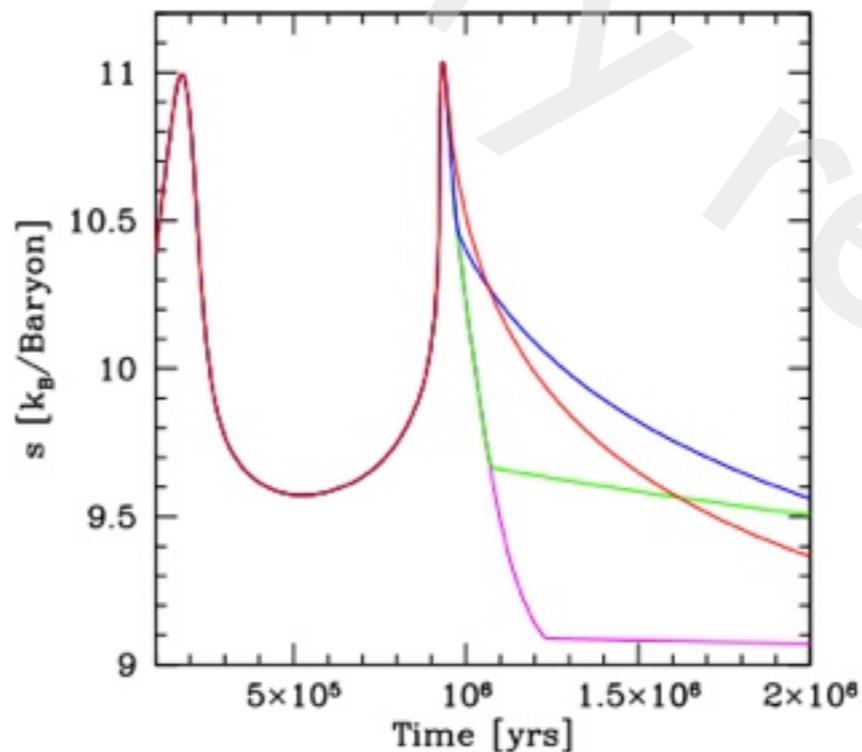
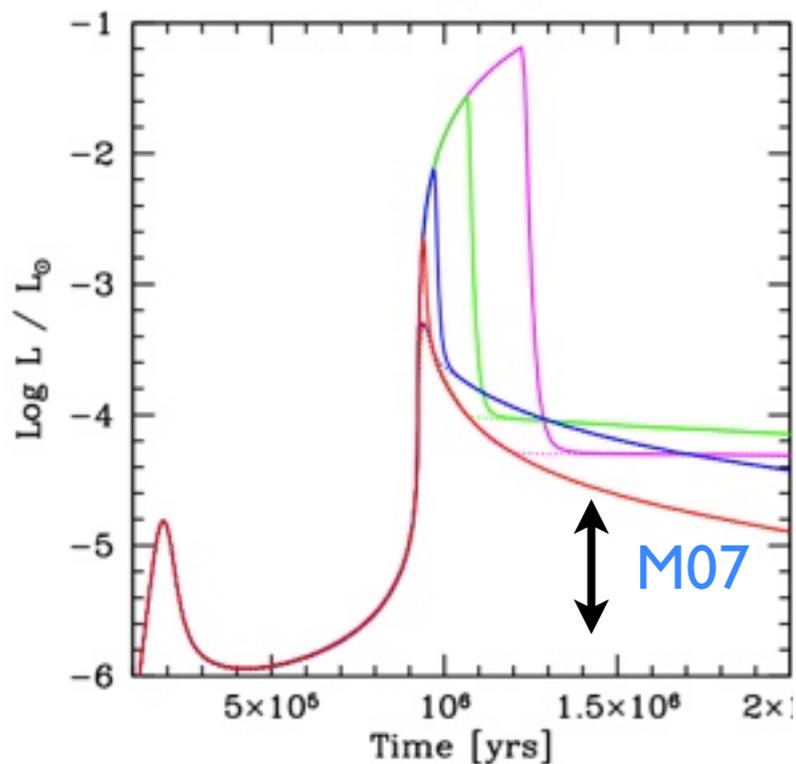
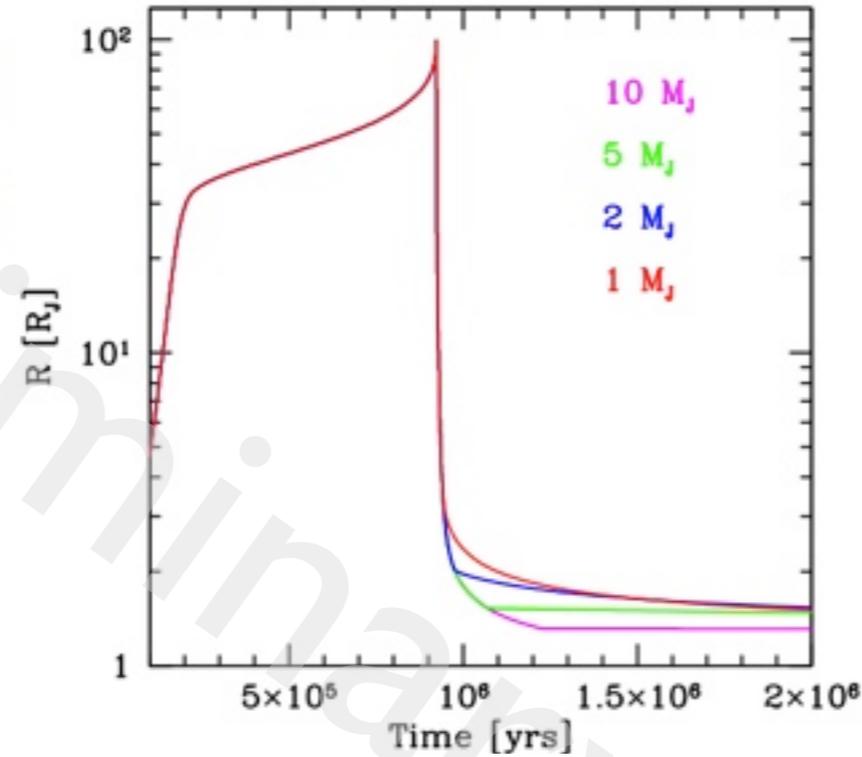
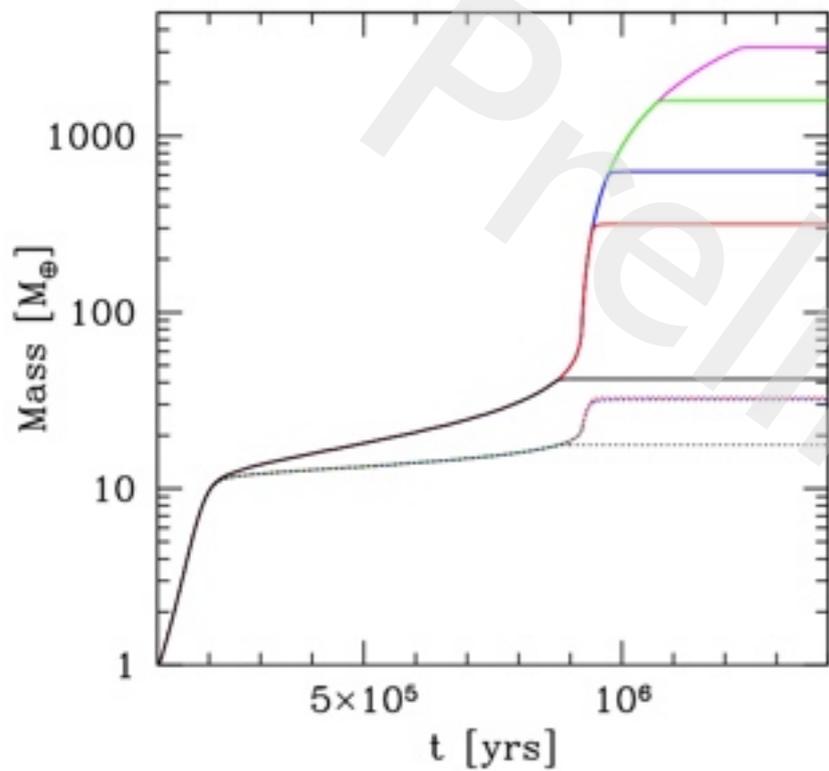
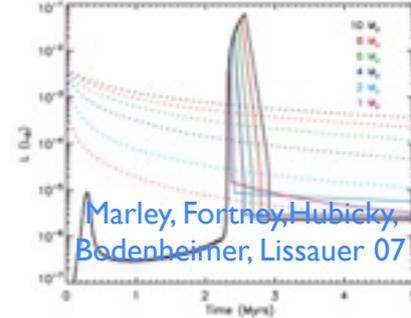
Use updated model to revisit luminosity of **young Jupiters**. Same model as normally used for population synthesis. Also as illustration that population synthesis formation model contains quite some **physics**.

- **Hot** start: Start with fully formed planet, at (arbitrarily) large luminosity (entropy) and radius. Outcome of direct collapse / disk instability (?).
- **Cold** start: Gradually build up planet (core accretion). Accrete gas through accretion shock. Radiative loss of gravitational potential energy liberated at the shock. Lower initial luminosity and radius.



Currently known direct imaging planets are better reproduced by hot start models.

New *cold* start models



- Initial conditions like **J1** in Pollack+ 1996. 1, 2, 5, 10 M_J . In situ. $\dot{M}_{\max}=0.01 M_e/\text{yr}$. Low grain opacity.

- **ALL** accretional luminosity radiated away at the shock.

- As in Marley+ 2007, **low** L at end of formation. But somewhat higher $L \Rightarrow$ see below.

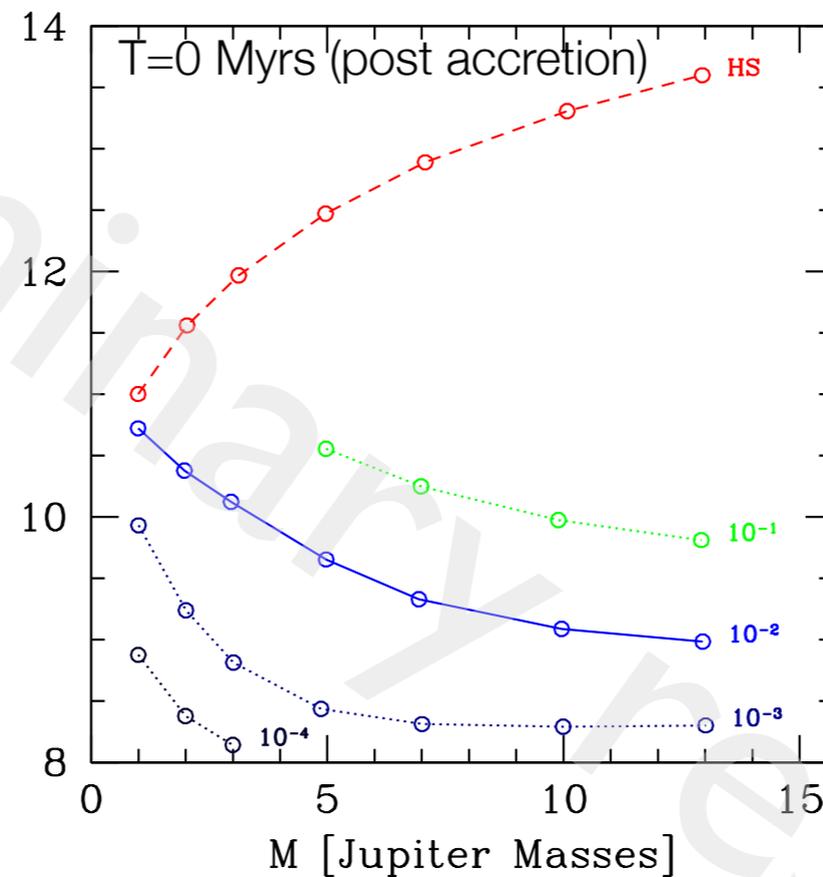
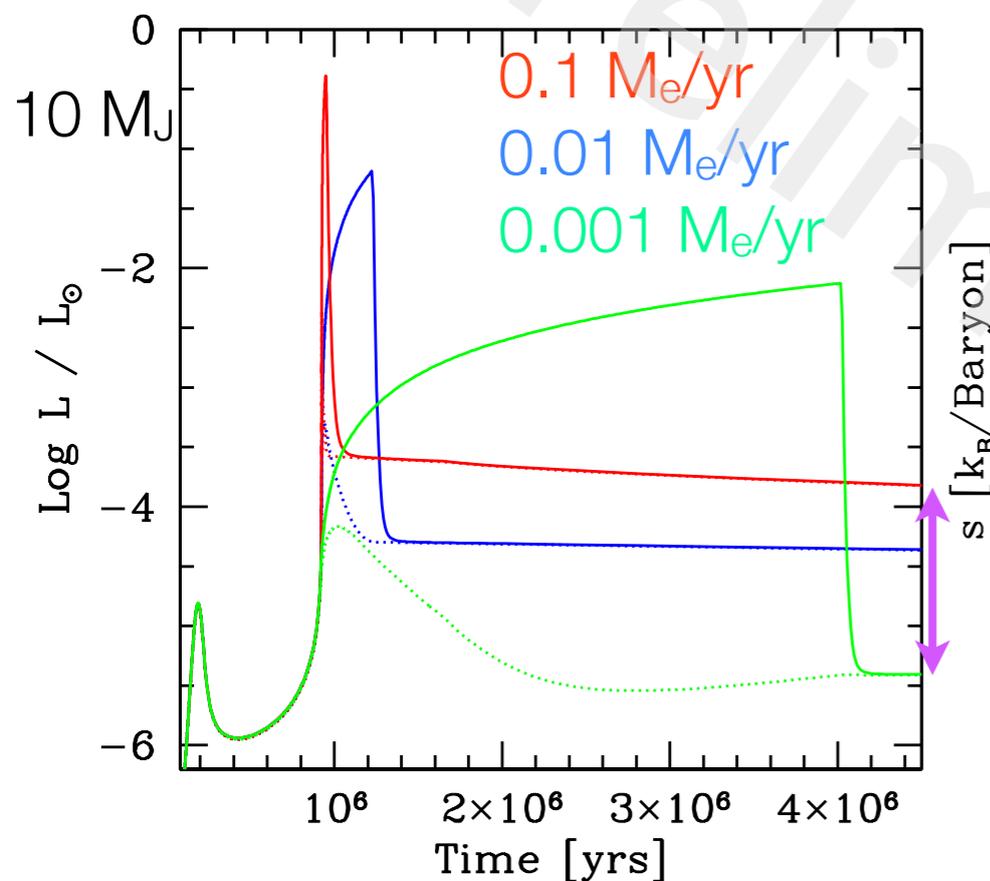
- Accretion of low entropy material.

*Unique?
Dependence on
parameters?*

Luminosity is a strong function of the gas accretion rate

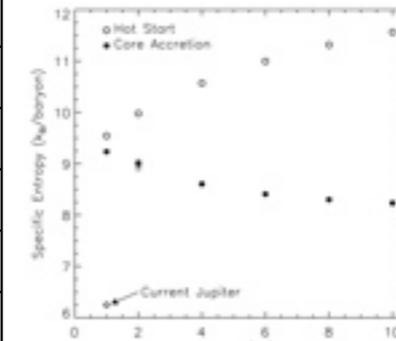
Gas accretion rate in runaway phase: given by disk (no more planet). Can vary!
 In full model, \dot{M}_{\max} calculated self consistently from disk model.

Here: Consider planet with final mass $10 M_J$, with \dot{M}_{\max} 0.1, 0.01 and 0.001 M_e/yr .

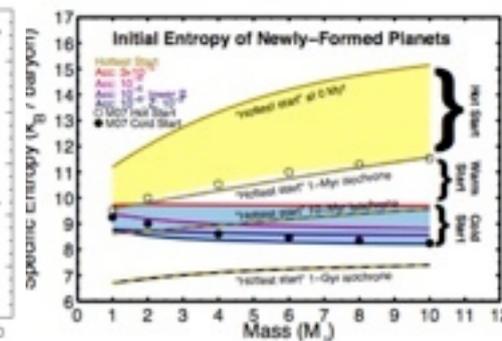


“Fork” diagram of initial specific entropies. High $s \Rightarrow$ high L .

Clear \dot{M}_{\max} dependence.



Marley+07

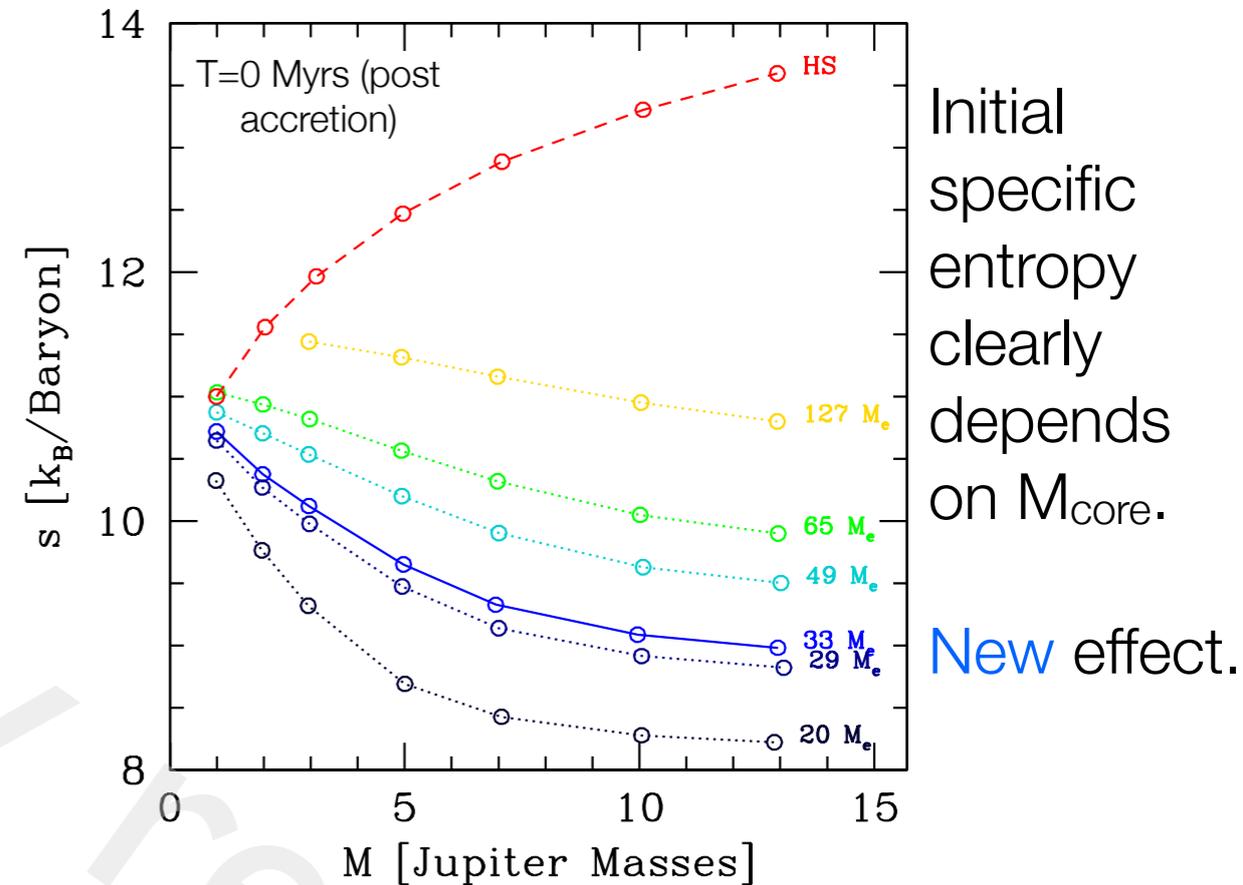
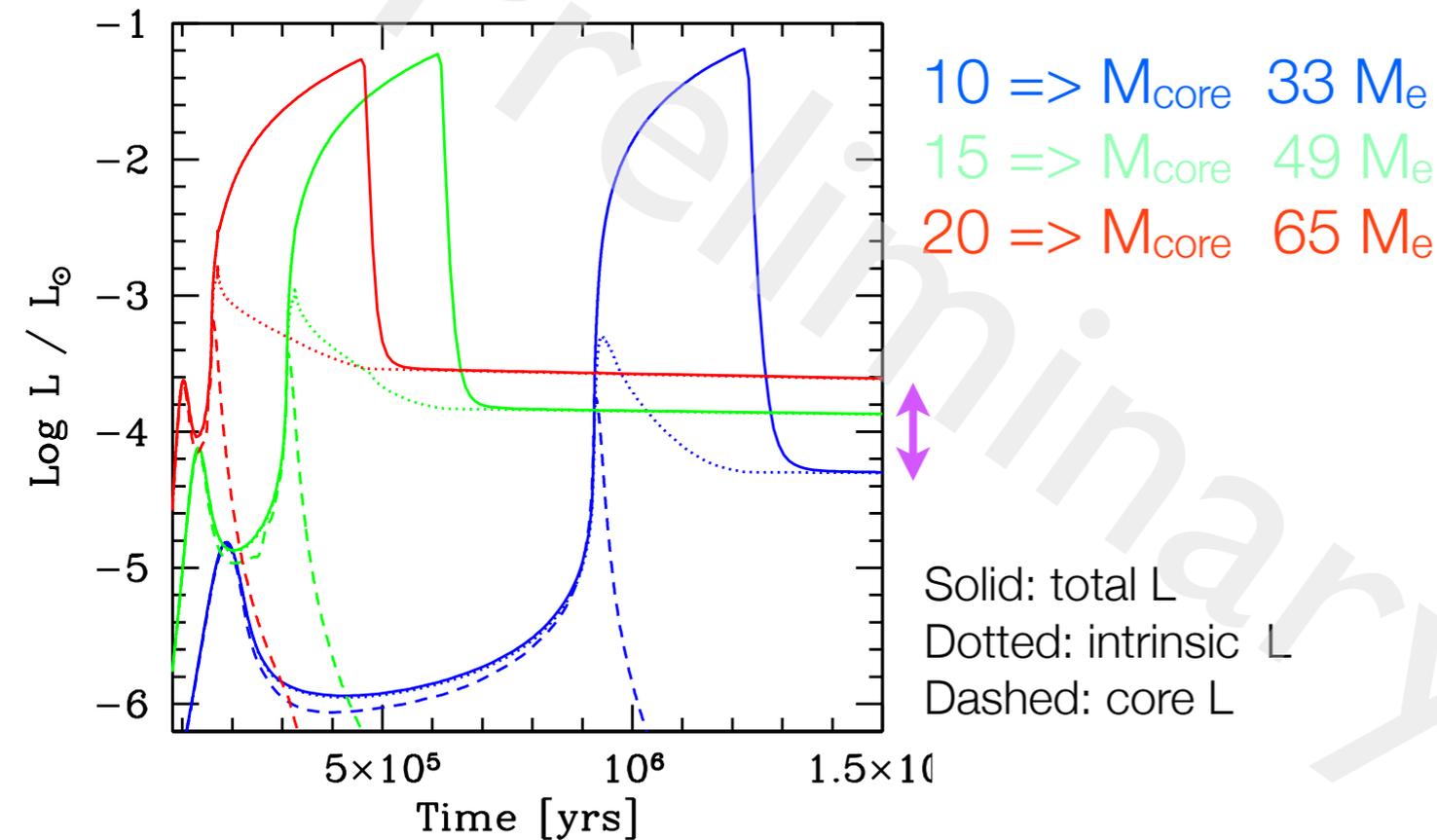


Spiegel & Burrows '11

- High gas accretion rate \Rightarrow high luminosity: almost 2 orders of magnitude difference.
- Accretion timescale vs. contraction timescale
- $t_{\text{accr}} \ll t_{\text{cont}}$ (high \dot{M}_{\max}): high fraction of mass already accreted while the planet is still big. Less low entropy matter accreted through the shock.
- The higher \dot{M}_{\max} , the closer we get to the hot start model which (formally) has a vanishing or very small t_{accr} .

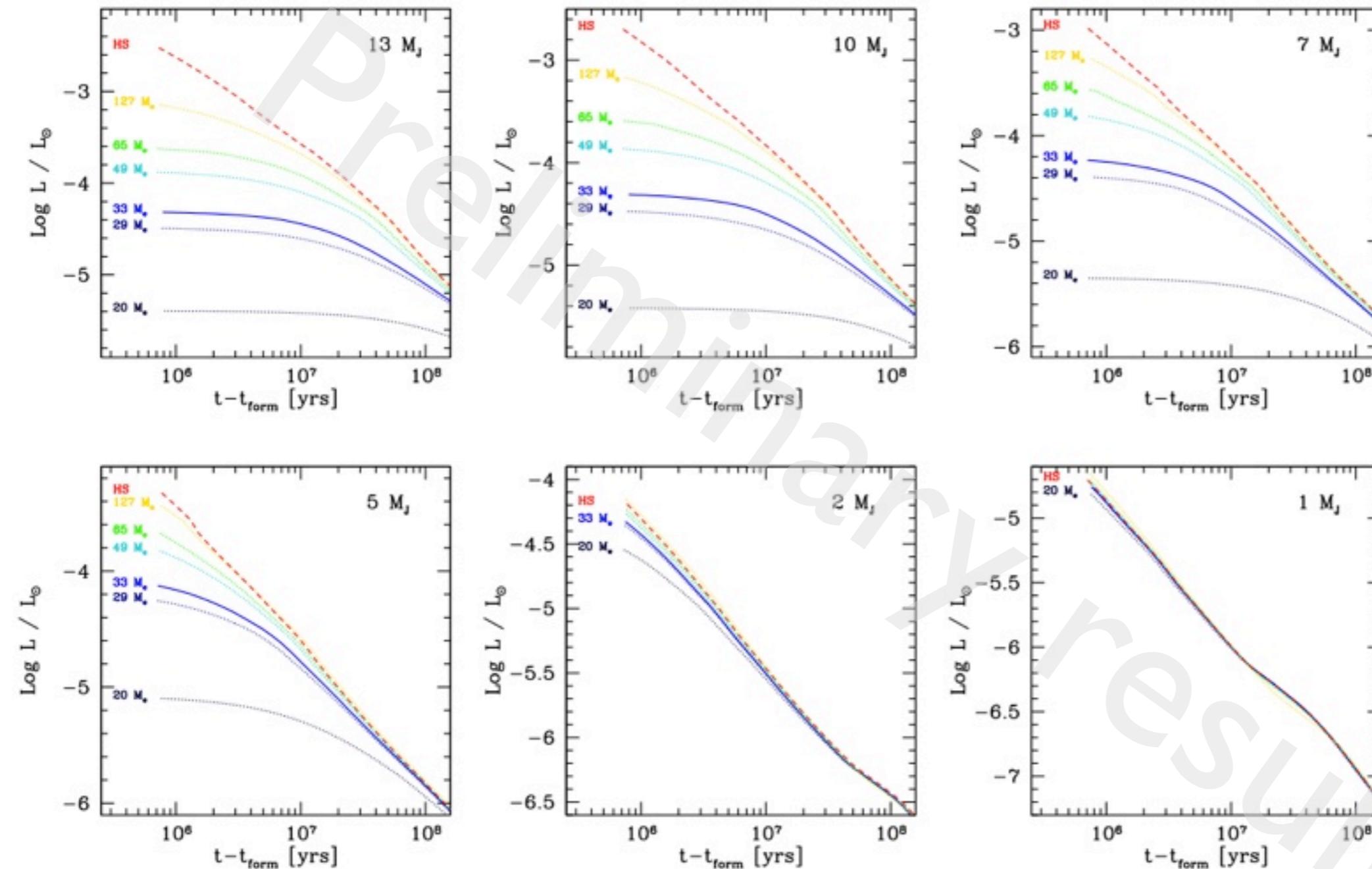
Luminosity is a function of the core mass /

Planetesimal surface density: The higher, the shorter time till runaway and the larger M_{core} .
 $10 M_J$, Initial planetesimal surface density $\Sigma_{p,0}$: 10, 15, 20 g/cm². $\dot{M}_{\text{max}} = 10^{-2} M_e/\text{yr}$



- During approach to runaway, L_{core} is a **non-negligible** contribution. Later decrease (R_{capt}).
- Matter contained in envelope when collapse starts is at **higher** entropy for higher L_{core} . Larger initial $R \Rightarrow$ weaker shock \Rightarrow less radiative loss \Rightarrow difference gets bigger. (**self sustaining process**). Different M_{core} explains also difference to Marley+2007.
- **Caveats**: depends on planetesimal random velocity & planetesimal-protoplanetary envelope interaction. No dependence of opacity on $\Sigma_{p,0}$ included here (cf Spiegel & Burrows 2011).

Luminosity is a strong function of the core mass II: $L(t, M, M_{\text{core}})$



Red line: Hot start

Other lines, cold start with different $\Sigma_{p,0}$ and thus M_{core}

Hot start and cold start with massive core identical after 1 to 10 Myrs.

Planets with $M_{\text{core}} > 100 M_e$ do exist (Leconte et al. 2010).

Very large differences in luminosity for higher mass planets depending on core mass, remaining up to 10^8 years. Again, there is no such thing as one unique “cold start” luminosity.

(Preliminary) Conclusion

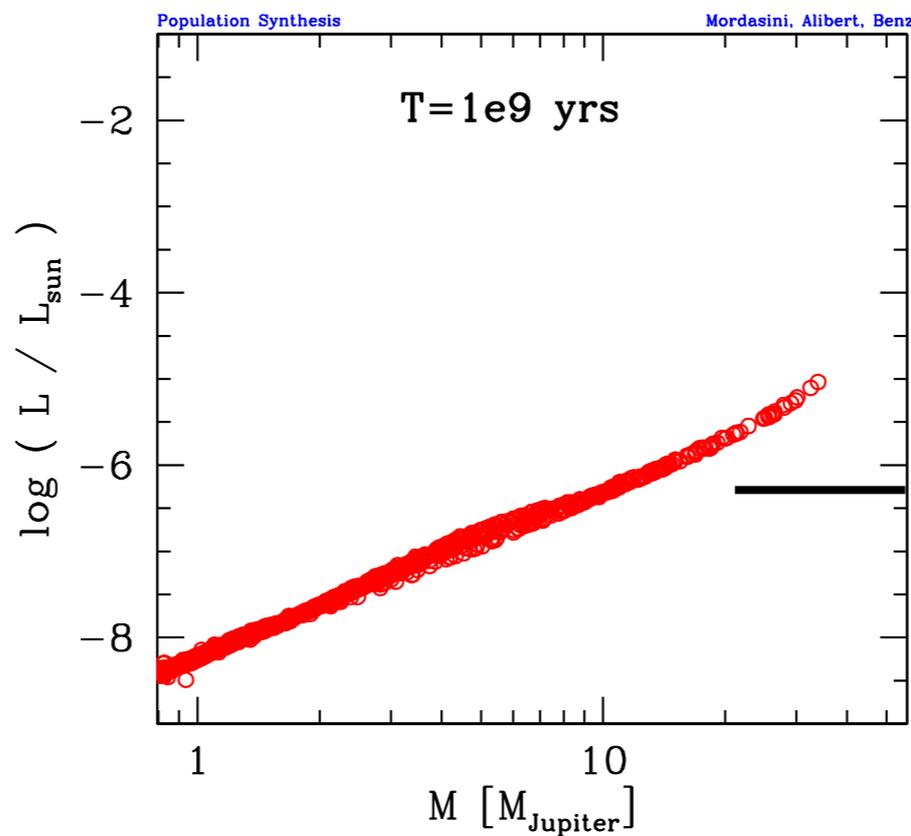
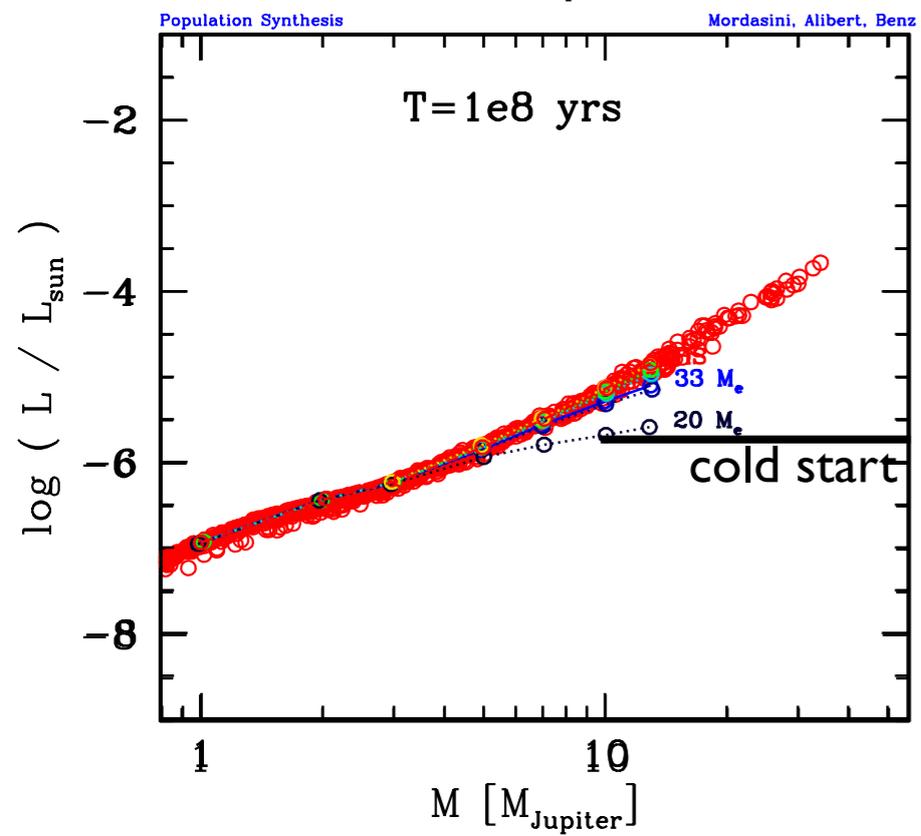
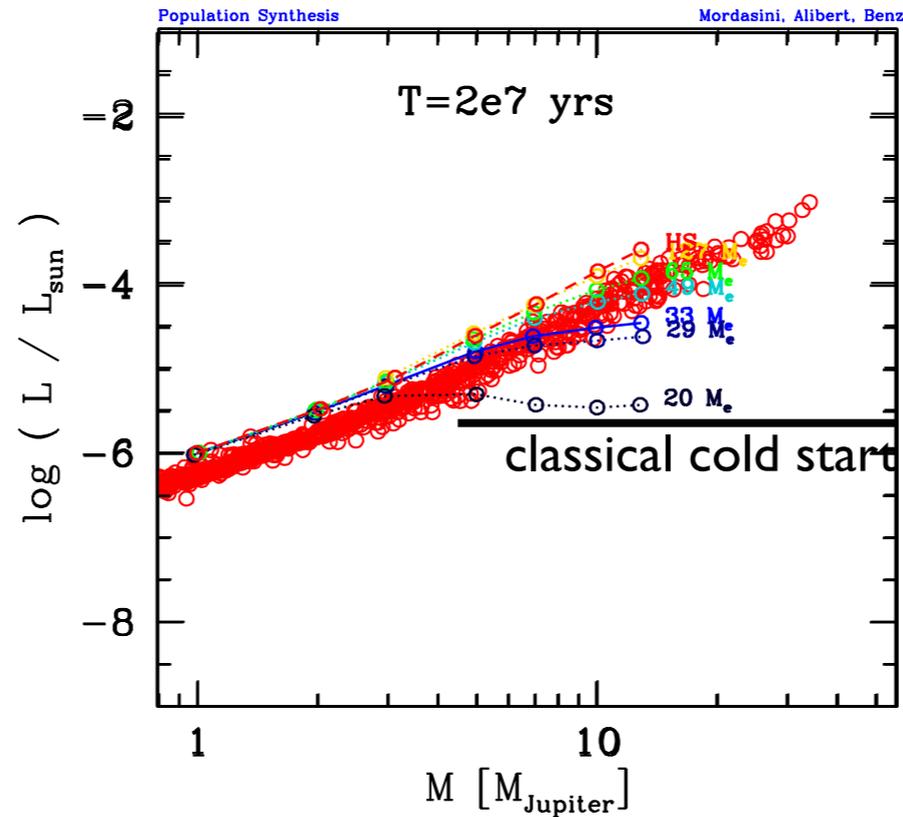
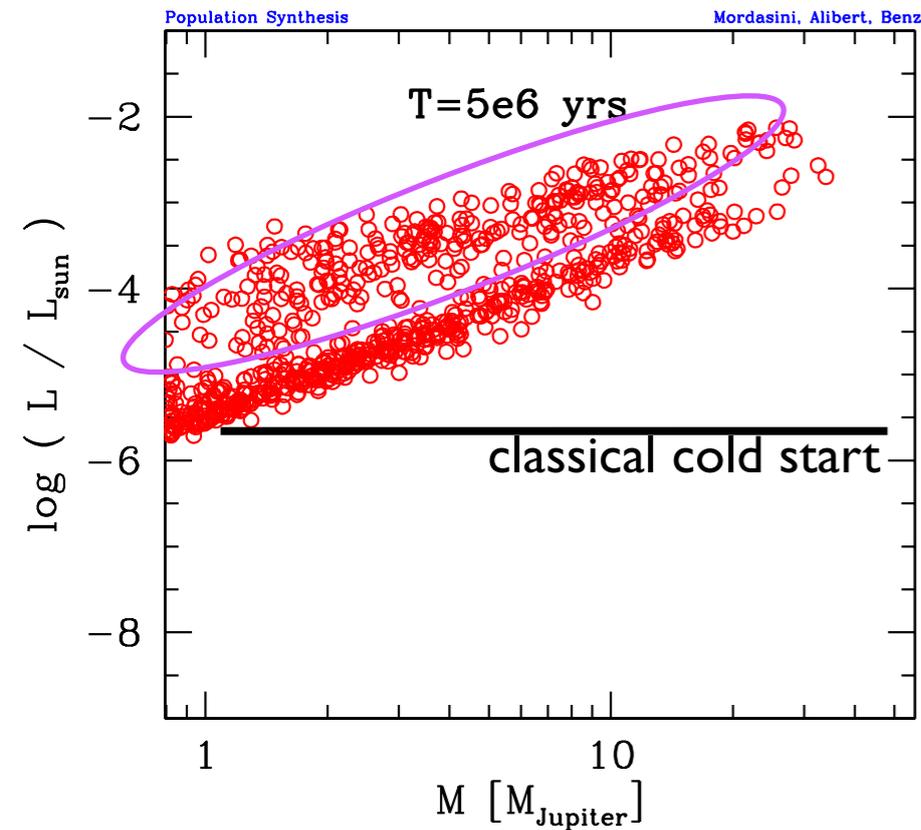
It is **impossible** to derive the **formation** mechanism or the **mass** from measuring the luminosity by **direct imaging**.

L can be anything.

Really?

Bring back in **population synthesis!**

Core Accretion is *hot*



Nominal population
 $M_{\text{star}}=1 M_{\odot}$
 $\alpha=7 \times 10^{-3}$
 $a > 0.3$ AU

Final conclusion from pop. synth.

Core accretion leads to L almost as *high* as Hot Start. $M_Z \geq 40 M_e$. At least from this point, CA reproduces also planets at large distances. It leads to a relatively *well* defined M - L relation.

Summary

- Presented upgraded core accretion model combining self-consistently **formation** and **evolution**.
- Allows **characterization** of planets in mass, semimajor axis, composition, radius and luminosity from tiny embryo to Gyr old planet.
- Updated type I migration rates allow to obtain populations with a radius distribution **similar** as observed by Kepler **without** scaling factors. Giant planets however get still too close.
- Core accretion leads to post formation **luminosities** almost as **high** as in the “hot start” scenario. No clear **difference** to direct collapse model. On the other hand, $L(M)$ is quite well defined.
- The model allows population synthesis results to be compared directly with **RV**, **transits** and **direct imaging** (and microlensing). This **combination** is the key to better **understanding** planet formation.